Green’s Functions Associated with One and Two Dimensional Problems

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Abstract: In mathematics a green’s function is type of function used to solve inhomogeneous differential equations subject to specific initial conditions or boundary conditions. Green’s functions provide an important tool when we study the boundary value problem. They also have intrinsic value for a mathematician. Also green’s functions in general are distribution, not necessarily proper function. Green functions are also useful for solving wave equation, diffusion equation and in quantum mechanics, where the green’s function of the Hamiltonian is a key concept, with important links to the concept of density of states.

Keywords: Green’s functions, inhomogeneous differential equations, boundary value problem.

1. Introduction

The green’s function as used in physics is usually defined with the opposite sign that is [1], [2],

\[ \ell G(x, s) = -\delta(x - s) \]

This definition does not change significantly any of the properties of the Green’s function in heat conduction we know that the Greens’ function represents that temperature at a field point due to a unit heat source applied at source point. In electro static the green’s function stand for the displacement in the solid due to the application of unit point force [3].

In this project construction of green’s function in one and two dimensions has shown. There are more than one way of constructing greens’ function (if it exists) but the result is always same. Due to this we can say that green’s function for a given linear system is unique. We start with the brief introduction of the Dirac delta or Dirac’s delta function which is not strictly a function in real sense of functions [4].

2. Dirac Delta Function

i. \( \delta(t-a) = 0 \) if \( t \neq 0 \)
ii. \( \delta(t-a)dt = 1 \)
iii. \( \delta(t-a)f(t)dt = f(a) \)
iv. \( \delta(t-a) = \delta(a-t) \)

3. Green’s Function Associated with One Dimensional Boundary Value Problem

Consider the following boundary value problem.

\[ M[y] = F(x) \rightarrow (1) \quad X_1 < x < X_2 \]

\[ B_1[y] = \infty \quad B_2[y] = \beta \]

M is defined by

\[ M = A_2(x) \frac{d^2}{dx^2} + A_1 \frac{d}{dx} + A_0(x) \]

From (1)

\[ A_2(x) \frac{d^2y}{dx^2} + A_1 \frac{dy}{dx} + A_0(x)y = F(x) \]

Divide by \( A_1(x) \) we will get

\[ \frac{d^2y}{dx^2} + \frac{A_1(x)}{A_2(x)} \frac{dy}{dx} + \frac{A_0(x)}{A_2(x)}y = \frac{F(x)}{A_2(x)} \rightarrow (2) \]

Let

\[ P(x) = e^{\int \frac{A_1(x)}{A_2(x)} dx} \]

Multiply equation (2) by \( P(x) \)

\[ P(x) \frac{d^2y}{dx^2} + \frac{A_1(x)}{A_2(x)} P(x) \frac{dy}{dx} + \frac{A_0(x)}{A_2(x)} P(x)y = \frac{F(x)P(x)}{A_2(x)} \]

\[ P(x)F(x) = \frac{d}{dx} \left[ P(x) \frac{dy}{dx} \right] + q(x)y = f(x) \]

\[ \ell[y] = f(x) \]

Where

\[ \ell = \frac{d}{dx} \left[ P(x) \frac{dy}{dx} \right] + q(x) \]

Consider the self-ad-joint boundary value problem

\[ \ell[y] = f(x) \]

\[ B_1[y] = \alpha, B_2[y] = \beta \rightarrow (3) \]

Now split the boundary value problem into the following boundary value problems.

\[ \ell[y] = 0 \]

\[ B_1[y] = \alpha, B_2[y] = \beta \rightarrow (4) \]

\[ \ell[y] = f(x) \]

\[ B_1[y] = 0, B_2[y] = 0 \rightarrow (5) \]

The solution of problem (3) is written as

\[ y = y_h + y_p \]
\( y_H \) is complementary function which satisfy the homogeneous differential equation
\[ \ell y_H = 0 \]
\( y_P \) is the particular solution to the inhomogeneous differential Equation
\[ \ell y_P = f(x) \]
The general solution of (4) is written as
\[ y_H = c_1 y_1 + c_2 y_2 \rightarrow (6) \]

Here \( y_1 \) & \( y_2 \) are solution of homogeneous Equation
\[ \ell [y] = 0 \quad \text{and} \quad c_1, c_2 \] are arbitrary constant and they can be determined by applying the boundary condition. Now consider the problem (5). Suppose that the solution of problem ‘5’ can be expressed in the integral form as
\[ y_P = -\int_{x_1}^{x_2} g(x, s)f(s)ds \rightarrow (7) \]

Where \( g(x, s) \) is the green’s function which is to be defined later. The negative sign in (7) describe the physical interpretation. Apply the differential operator \( \ell \) on both sides of \( \ell y_P = f(x) \)
\[ \ell [y_P] = \ell \left(-\int_{x_1}^{x_2} g(x, s)f(s)ds\right) \]

Since \( \ell \) & the integral operator commute each other.
\[ \ell [y_P] = -\int_{x_1}^{x_2} \ell (g)f(s)ds \rightarrow (i) \]

Also
\[ -\int_{x_1}^{x_2} \delta(x-g)f(s)ds = f(x) \rightarrow (ii) \]
And also we have
\[ \ell [y_P] = f(x) \rightarrow (iii) \]
From i, ii, and iii.
\[ -\int_{x_1}^{x_2} \ell (g)f(s)ds = -\int_{x_1}^{x_2} \delta(x, s)f(s)ds \]
\[ \Rightarrow -\int_{x_1}^{x_2} [\ell (g) + \delta(x-s)]f(s)ds = 0 \]
Since \( f(s) \) is arbitrary \( f(s) \neq 0 \)
\[ \ell (g) + \delta(x-s) = 0 \]
\[ \ell (g) = -\delta(x-s) \rightarrow (8) \]
\( \delta(x-s) \) is dirac delta function.

To determine the unique green’s function condition (8) is not enough. We have to determine other condition also from homogeneous boundary condition.
\[ B_1[y] = 0 \quad , \quad B_2[y] = 0 \]

And
\[ B_1[y_P] = 0 \quad , \quad B_2[y_P] = 0 \]

Since \( f(s) \) Can be almost any function the above relations are satisfied only if.
\[ B_1[g] = 0 \quad , \quad B_2[g] = 0 \rightarrow (9) \]
Hence the green’s function we are looking for is solution of the following boundary problem
\[ \ell [g] = -\delta(x-s) \rightarrow (10) \]
Related with boundary condition
\[ B_1[g] = 0 \quad , \quad B_2[g] = 0 \]
Where’s is a fix value lies between \( x_1 \) and \( x_2 \) And. The above problem is similar to that given equation in (5) only the forcing function in (9) is delta function rather than arbitrary function \( f(x) \). This means that solving the problem for \( g \) is simpler then solving the corresponding problem ‘\( y \)’.

And once the green’s function has been determined for particular operator \( \ell \) and set of boundary conditions it may be used for solving problem (5) for any number of time where only the function \( f(x) \) changes from problem to problem. It is this feature of green’s function that makes it most useful in application [5]. The green’s function \( g(x,s) \) associated with boundary value problem
\[ \ell [y] = f(x) \]
\[ B_1[y] = \alpha \quad , \quad B_2[y] = \beta \]
Where
\[ \ell = d \frac{d}{dx} \left[ P(x) \frac{d}{dx} \right] + q(x) \]
Satisfying………………
(a) \[ \ell [g] = -\delta(x-s) \quad x_1 < x < x_2 \]
(b) \[ B_1[g] = 0 \quad , \quad B_2[g] = 0 \]
(c) \[ g(s^+,s) = g(s^-,s) \]
(d) \[ \frac{\partial g(s^+,s)}{\partial x} - \frac{\partial g(s^-,s)}{\partial x} = -\frac{1}{P(s)} \]
The condition ‘\( \delta \)’ is called jumped discontinuity of green’s function at \( x=5 \).

Based on above condition an exploit formula for the green’s function can be designed. It is observed from the condition
(a) if either \( x<s \) or \( x \geq s \) then
By definition of Dirac Delta function if 
\\[g(x) = 0\] 
By definition of Dirac Delta function if 
\\[g(x) = 0\] 
\\[\\begin{align*}
\ell[g] = 0 & \text{ By definition of Dirac Delta function if } \\
& \text{ such that } B_1[Z_1] = 0 \quad B_2[Z_2] = 0 \\
& \text{ so, from condition (a) and (b) the green’s function has following form} \\
g(x,s) = \begin{cases} 
\frac{u(s)Z_1(x)}{s} & x < s \\
\frac{v(s)Z_2(x)}{s} & x > s
\end{cases}
\end{align*}\\]
where u and v are function to be determined imposing the condition c and d. the unknown function u & v must be chosen such that 
\\[v(s)Z_2(x) - u(s)Z_1(x) = 0\\]
\\[v(s)Z_2'(x) - u(s)Z_1'(x) = -\frac{1}{P(s)}\\]
\\[\text{By solving the above two simultaneous equation we will get}\\]
\\[v(s) = \frac{-Z_1(x)}{P(s)W(Z_1Z_2)(s)}\\]
\\[W(Z_1Z_2) = Z_1Z_2 - Z_2Z_1\\]
Since Z_1 and Z_2 are such solution we can write
\\[P(s)W(Z_1Z_2)(s) = P(s)W(Z_1Z_2)(x)(x) = c\\]
The green’s function can be written as
\\[g(x,s) = \begin{cases} 
\frac{-Z_1(x)Z_2(s)}{c} & x < s \\
\frac{-Z_1(x)Z_2(s)}{c} & s \leq x < x_2
\end{cases}\\]
Green’s function is symmetric in x and s, i.e. g(x,s) = g(s,x)

1) Green’s function associated with two dimensional problems:
To introduce green’s function in two dimensions, we consider the PDE (Poisson’s equation which here represents static deflection of a rectangular membrane).
\\[\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = f(x,y) \rightarrow (11)\\]
Here f(x,y) represent the external load per unit area, divided by T(tension in the membrane, which here has the dimension of force per length)
The B.C.S. in this case are 
\\[u(0,y) = u(a,y) = 0, u(x,0) = u(x,b) = 0 \rightarrow (12)\\]
A concentrated force acting at a point (x’,y’) may be simulated by the two dimension delta function
\\[\frac{F}{\tau} \delta(x-x')\delta(y-y')\\]
Let 
\\[G(x',y'|y') = 0\\]
be the green’s function associated with the problem
\\[\frac{\partial^2 G}{\partial x^2} + \frac{\partial^2 G}{\partial y^2} = \frac{F}{\tau} \delta(x-x')\delta(y-y') \rightarrow (13)\\]
\\[G(0,y) = G(a,y) = 0, G(x,0) = G(x,b) = 0 \rightarrow (14)\\]
Then the solution of the problem (11) and (12) is given by
\\[u(x,y) = \iint_G G(x',y'|y')f(x',y')dx'dy'\\]
We first find a complete set of Eigen functions of the associated homogeneous Eigen value problem viz.
\\[\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right)u_{ij}(x,y) = \lambda u_{ij}(x,y) \text{ with the same B.C.S as on } u(x,y).\\]
We find
\\[\lambda_{mn} = \left(\frac{m^2\pi^2}{a^2} + \frac{n^2\pi^2}{b^2}\right)\\]
\\[\text{And}\\]
\\[u(x,y) \equiv u_{mn}(x,y) = \frac{2}{\sqrt{ab}} \sum_{m=1,n=1}^{\infty} A_{mn}(x',y') \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b}\\]
Which have been normalized to unity. Now the required green’s function can be represented as
\\[G(x',y'|y') = \frac{2}{\sqrt{ab}} \sum_{m=1,n=1}^{\infty} A_{mn}(x',y') \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b}\\]
Now the required green’s function can be represented as
\\[G(x',y'|y') = \frac{2}{\sqrt{ab}} \sum_{m=1,n=1}^{\infty} A_{mn}(x',y') \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b}\\]
Substituting this into (13)
\\[\sum_{n=1}^{\infty} \left[ -\frac{m^2\pi^2}{a^2} - \frac{n^2\pi^2}{b^2} \right] A_{mn}(x',y') \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b} = \frac{F}{\tau} \delta(x-x')\delta(y-y')\\]
Multiplying by \(\sin(m\pi x/a)\sin(n\pi y/b)\) and integrating w.r.t x and y we have.
\\[-\frac{2}{\sqrt{ab}} \left( \frac{m^2\pi^2}{a^2} + \frac{n^2\pi^2}{b^2} \right) A_{mn}(x',y') = \frac{F}{\tau} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b}\\]
Or
\\[-\left( \frac{m^2\pi^2}{a^2} + \frac{n^2\pi^2}{b^2} \right) A_{nn}(x',y') = -\frac{2}{\sqrt{ab}} \frac{F}{\tau} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b}\\]
Hence on substitution
\\[G(x',y'|y') = \frac{4}{\sqrt{ab}} \sum_{m,n}^{\infty} \left( \frac{a^2b^2}{m^2\pi^2 + a^2n^2\pi^2} \right) \frac{F}{\tau} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b} \sin \frac{m\pi x}{a} \sin \frac{n\pi y}{b}\\]
4. Conclusion

This paper provides the introduction of green’s function associated with one and two dimensional problems to solve inhomogeneous differential equations subject to specific initial conditions or boundary conditions. We start with the brief introduction of the Dirac delta or Dirac’s delta function and then explore Green’s Function Associated with one dimensional boundary value problem. Finally green’s function in two dimensions, we consider the PDE’s e.g. Poisson’s equation which here represents static deflection of a rectangular membrane.

References